

NUMERICALLY INVARIANT SIGNATURE CURVES

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ABSTRACT. Corrected versions of the numerically invariant expressions for the affine and Euclidean signature of a planar curve given in [1] are presented. The new formulas are valid for fine but otherwise arbitrary partitions of the curve. We also give numerically invariant expressions for the four differential invariants parametrizing the three dimensional version of the Euclidean signature curve, namely the curvature, the torsion and their derivatives with respect to arc length.

1. INTRODUCTION

The concept of signature was introduced by Calabi et al. [1] as “a new paradigm for the invariant recognition of visual objects.” For the case of planar curves, the signature is defined as follows: if $\alpha(s)$ is a smooth curve in \mathbb{R}^2 parametrized by arc length s and if G is any finite dimensional transformation group acting transitively on \mathbb{R}^2 , then the signature curve S of $\alpha(s)$ with respect to G is given parametrically as $(\kappa(s), \kappa_s(s))$, where κ is the G -invariant curvature and κ_s its derivative with respect to arc length. One of the consequences of a theorem proved by Cartan [3] is that S fully determines the given curve α modulo G , provided that κ is never zero.

The signature curve can therefore be used to program a computer for recognizing curves modulo certain group transformations. However differential invariants are high order derivatives and hence very sensitive to round off errors and noise. The idea of writing a numerical scheme in terms of joint invariants was also introduced by Calabi et al. [1]. Hoping to obtain less sensitive approximations, they suggested to find numerical expressions for κ and κ_s in terms of joint invariants. A joint invariant of the action of a group on a manifold is a real valued function J which depends on a finite number of points x_1, \dots, x_n of the manifold and which remains unchanged under the simultaneous action of the group G on the point configuration, i.e. $J(x_1, \dots, x_n) = J(g \cdot x_1, \dots, g \cdot x_n) \forall g \in G$. For example, the Euclidean distance is a joint invariant of the action of the Euclidean group on \mathbb{R}^2 . Expressing differential invariants in terms of joints invariants results in a G -invariant numerical approximation. In their paper, Calabi et al. proposed numerically invariant expressions for κ and κ_s for two specific group actions, namely the proper Euclidean group and the equi-affine group.

But contrary to their claim, the expressions given for κ_s are not convergent for arbitrary partitions of the curve. In the next section, we give correct formulas for approximating κ_s and explain why the old ones do not work in general. In order to prove our claims, we also compute the resulting numerical signatures in a practical example and with different partitions. We then compare with the signatures obtained with the old formulas.

Cartan's theorem also provides us with a way to characterize curves in \mathbb{R}^3 modulo group transformation. The generalization of the signature curve is a curve in \mathbb{R}^4 determined by four differential invariants: the G -invariant curvature κ , its derivative κ_s with respect to G -invariant arc length s , the G -invariant torsion τ and its derivative τ_s with respect to s . For practical applications, we are interested in the case where G is the proper Euclidean group. Following the example of [1], we have found approximations for κ , κ_s , τ and τ_s in terms of the simplest joint invariant of the action of the Euclidean group on \mathbb{R}^3 : the Euclidean distance. The results are given in section 3. That section also contains the results of numerical tests performed on a space curve with different parametrizations.

2. CORRECTIONS FOR THE CASE OF A PLANAR CURVE

In principle, one must keep track of the order of the approximation when manipulating an approximation. This is especially true when trying to approximate a derivative using an approximate expression. For example, if $\tilde{c}(t) = c(t) + O(t)$ is a first order approximation for $c(t)$, then for $0 < h, t \ll 1$

$$\frac{\tilde{c}(t) - \tilde{c}(0)}{t} = \frac{c(t) + O(t) - c(0)}{t} \approx c'(0) + O(t^0)$$

is in general *not* an approximation for $c'(0)$.

By keeping track of the order of approximation throughout the computations performed in [1], one finds out that the expressions given for κ_s will converge only for a 2^{nd} order approximation of κ . In the case of a regular partition of the curve, the approximations used for the curvature are of second order, as one can see from equation (1) below. This fortunate fact explains the correct numerical results presented in [1]. However for a generic partition there is no guarantee of success and failure is likely as we show by example.

In the following we present the corrected formulas together with their justification for the Euclidean and affine group action..

2.1. Euclidean Group Action. Let P_{i-2} , P_{i-1} , P_i , P_{i+1} , and P_{i+2} be 5 consecutive points on a smooth planar curve. Denote by $d(P_j, P_k)$ the Euclidean distance between the points P_j and P_k . As illustrated in figure 1, let

$$\begin{aligned} a &:= d(P_{i-1}, P_i), & c &:= d(P_{i-1}, P_{i+1}), & e &:= d(P_i, P_{i+2}), & g &:= d(P_{i-2}, P_{i-1}) \\ b &:= d(P_i, P_{i+1}), & d &:= d(P_{i+1}, P_{i+2}), & f &:= d(P_{i-1}, P_{i+2}) \end{aligned}$$

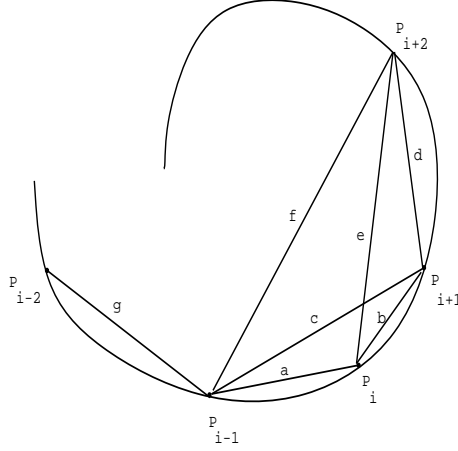
It was shown in [1] that $\tilde{\kappa}(P_i) = \pm \frac{4\Delta}{abc}$, where Δ denotes the area of the triangle with sides of length a , b and c , is a good approximation for the Euclidean curvature at P_i . In fact, the following expansion has been proven to be valid for small a and b .

$$(1) \quad \tilde{\kappa}(P_i) = \kappa(P_i) + \frac{1}{3}(b-a)\kappa_s(P_i) + O(2).$$

Observe that $\tilde{\kappa}$ depends only on the distances between the three points. Since the distance is an Euclidean invariant, we thus have a Euclidean numerically invariant expression for the curvature.

When $a \neq b$, the approximation obtained is only of first order. Therefore care must be taken when approximating the derivative $\kappa_s(P_i)$. The formulas used in [1]

FIGURE 1.



are

$$(2) \quad \tilde{\kappa}_{s,1}(P_i) := \frac{\tilde{\kappa}(P_{i+1}) - \tilde{\kappa}(P_i)}{b},$$

$$(3) \quad \tilde{\kappa}_{s,1}(P_i) := \frac{\tilde{\kappa}(P_{i+1}) - \tilde{\kappa}(P_{i-1})}{c}.$$

We would like to have expressions that do not assume a priori conditions on the parametrizations. Hence we propose the following numerically Euclidean invariant expression for $\kappa_s(B)$:

$$(4) \quad \tilde{\kappa}_{s,3}(P_i) := 3 \cdot \frac{\tilde{\kappa}(P_{i+1}) - \tilde{\kappa}(P_i)}{a + b + d},$$

$$(5) \quad \tilde{\kappa}_{s,4}(P_i) := \frac{3}{2} \cdot \frac{\tilde{\kappa}(P_{i+1}) - \tilde{\kappa}(P_i)}{a + b + d} + \frac{3}{2} \cdot \frac{\tilde{\kappa}(P_i) - \tilde{\kappa}(P_{i-1})}{a + b + g},$$

$$(6) \quad \tilde{\kappa}_{s,5}(P_i) := 3 \cdot \frac{\tilde{\kappa}(P_{i+1}) - \tilde{\kappa}(P_{i-1})}{2a + 2b + d + g}.$$

Claim 2.1. *Formulas (4), (5) and (6) converge to $\kappa_s(P_i)$ as $a, b, d, g \rightarrow 0$*

Proof. By (1) we have

$$\tilde{\kappa}(P_i) = \kappa(P_i) + \frac{1}{3}(b - a)\kappa_s(P_i) + O(2)$$

and

$$\tilde{\kappa}(P_{i+1}) = \kappa(P_{i+1}) + \frac{1}{3}(d - b)\kappa_s(P_{i+1}) + O(2).$$

Assume that $\alpha = \alpha(s)$ is a smooth curve parametrized by arc length, that $P_i = \alpha(0)$ and that $P_{i+1} = \alpha(h)$. We can write

$$\kappa(P_{i+1}) = \kappa(P_i) + h\kappa_s(P_i) + O(2) \approx \kappa(P_i) + b\kappa_s(P_i) + O(2)$$

and

$$\kappa_s(P_{i+1}) = \kappa_s(P_i) + O(1).$$

Thus

$$\tilde{\kappa}(P_{i+1}) - \tilde{\kappa}(P_i) = \frac{1}{3}(d - 2b + a)\kappa_s(P_i) + b\kappa_s(P_i) + O(2).$$

Rearranging, we obtain

$$3 \cdot \frac{\tilde{\kappa}(P_{i+1}) - \tilde{\kappa}(P_i)}{a + b + d} = \kappa_s(P_i) + O(1)$$

Therefore (4) is valid.

Equation (5) is just a symmetric formula obtained by averaging (4) with its reflected version. We obtain (6) by expanding $\tilde{\kappa}(P_{i+1}) - \tilde{\kappa}(P_{i-1})$ about P_i in a similar way. \square

Figures 2-8 show that our formulas are just as good as (2) and (3) when the parametrization is very regular. But when the parametrization is not, we see that (2) and (3) diverge from the exact solution, while (4), (5) and (6) remain closer to it. Thus we believe that our formulas are better for a generic partition of α . Observe that formula (4) suffers from a small bias due to its asymmetry. This phenomenon was also observed for formula (2) in [1].

2.2. Affine Group Action. Let α be a convex smooth curve. Let P_0, P_1, P_2, P_3, P_4 be five consecutive points on α . Denote by $[ijkl]$ the signed area of the parallelogram whose sides are $P_i - P_j$ and $P_k - P_l$ and by $[ijk]$ the signed area of the parallelogram whose sides are $P_i - P_j$ and $P_i - P_k$. Define T and S at P_2 by

$$4 \cdot T(P_2) = \prod_{0 \leq l < m < n \leq 4} [lmn],$$

$$\begin{aligned} 4 \cdot S(P_2) &= [013]^2[024]^2[1234]^2 + [012]^2[034]^2[1324]^2 \\ &\quad - 2[012][034][013][024]([123][234] + [124][134]). \end{aligned}$$

Observe that the expressions for T and S are equi-affine invariant. The following is a numerically invariant approximation for the affine curvature at P_i :

$$(7) \quad \tilde{\kappa}(P_i) = \frac{S(P_i)}{T(P_i)^{2/3}}.$$

Indeed the following expansion has been shown to be valid for sufficiently close points [1]

$$(8) \quad \tilde{\kappa}_i = \kappa + \frac{1}{5} \left(\sum_{j=i-2}^{i+2} L_j \right) \frac{d\kappa}{ds} + \dots$$

where L_j denotes the signed affine arc length of the conic from P_i to P_j and the higher order terms are quadratic in the L_j 's. Again, since this is a first order approximation, care must be taken when approximating the derivative of κ with respect to affine arc length s .

The formula used in [1] is

$$(9) \quad \tilde{\kappa}_s = \frac{\tilde{\kappa}(P_{i+1}) - \tilde{\kappa}(P_{i-1})}{\tilde{S}_i + \tilde{S}_{i-1}},$$

where \tilde{S}_j is the triangular approximation to the affine arc length from P_j to P_{j+1} . This formula would be valid only if $\tilde{\kappa}$ were a second order approximation of κ . However, this is not the case unless $\sum_{j=i-2}^{i+2} L_j$ is small compared to all of the L_j 's. We would like a more general expression. Hence we propose the following numerically affine invariant expression for κ_s which does not assume a priori conditions on the parametrizations:

$$(10) \quad \tilde{\kappa}_s^1 = 5 \cdot \frac{\tilde{\kappa}(P_{i+1}) - \tilde{\kappa}(P_{i-1})}{\tilde{S}_{i-3} + 2\tilde{S}_{i-2} + 2\tilde{S}_{i-1} + 2\tilde{S}_i + 2\tilde{S}_{i+1} + \tilde{S}_{i+2}}$$

The convergence can be proved by an argument similar to the one used in the Euclidean case.

Figures 9-12 show the results obtained with our formula using different partitions of a curve. Comparison is made with the signature curve obtained with (9). Just like in the Euclidean case, we see that equations (9) and (10) seem just as good when the partition of the curve is regular. But when the partition is not regular, (9) diverges from the exact solution while (10) remains valid.

3. NUMERICALLY INVARIANT EUCLIDEAN SIGNATURE FOR A SPACE CURVE

The signature curve S of a space curve is given by $S = (\kappa, \kappa_s, \tau, \tau_s)$. In the following section, we present some numerically Euclidean invariant approximations for S for the case of the Euclidean group action on \mathbb{R}^3 . This is achieved by expressing each of the differential invariants constituting S in terms of Euclidean distances, which is the simplest joint invariant of the Euclidean group action on \mathbb{R}^3 .

3.1. Numerically Invariant Expressions for κ and κ_s . Let P_{i-1} , P_i and P_{i+1} be three consecutive points on a space curve and let a , b and c be their mutual distances as illustrated in figure 1. Denote by Δ the area of the triangle with sides a , b and c . The curvature of a space curve is defined exactly the same way as the curvature of a planar curve: it is the inverse of the radius of the osculating circle at p . So we expect that $\tilde{\kappa} = \pm \frac{4\Delta}{abc}$ is a good approximation for κ at P_i . Indeed, a Mathematica routine using the local canonical form of a curve in \mathbb{R}^3 gave us the following Taylor series expansion.

$$\begin{aligned} \tilde{\kappa} &= \kappa + \frac{(b-a)}{3}\kappa_s + (a^2 + b^2)\left(\frac{\kappa_{ss}}{12} - \frac{\kappa\tau^2}{36}\right) - ab\left(\frac{\kappa_{ss}}{12} + \frac{\kappa\tau^2}{36}\right) \\ &+ (b^3 - a^3)\left(\frac{\kappa^2\kappa_s}{40} - \frac{7\kappa_s\tau^2}{570} - \frac{\kappa\tau\tau_s}{45} + \frac{\kappa_{sss}}{60}\right) \\ &+ (a^2b - ab^2)\left(-\frac{\kappa^2\kappa_s}{60} + \frac{\kappa_s\tau^2}{180} + \frac{\kappa\tau\tau_s}{180} + \frac{\kappa_{sss}}{60}\right) + \dots \end{aligned}$$

Note that if we set $\tau = 0$, then our result agrees with the 2-dimensional case. Also observe that τ does not appear in the first two terms of the expansion. Therefore the expressions for κ_s given in the planar case ((4), (5) and (6)) are valid numerically invariant expressions for κ_s in the three dimensional case.

3.2. Numerically Invariant Expressions for τ and τ_s . The Euclidean invariant torsion τ is defined as the derivative of the angle of the osculating plane with respect to arc-length. Let $-1 \ll -\delta < 0 < \epsilon_1 < \epsilon_2 \ll 1$. Suppose $P_{i-1} = \alpha(-\delta)$, $P_i = \alpha(0)$, $P_{i+1} = \alpha(\epsilon_1)$ and $P_{i+2} = \alpha(\epsilon_2)$ are four consecutive points on a space curve α . Let a, b, c, d, e and f be their mutual distances as illustrated in figure 1. Denote by H the height of the tetrahedron with sides a, b, c, d, e and f with respect to P_{i+2} and write $\Delta_{l_1 l_2 l_3}$ for the area of a triangle with sides of length l_1, l_2 and l_3 .

We propose the two following numerically invariant expressions for the torsion at P_i :

$$(11) \quad \tilde{\tau}_1 = 6 \cdot \frac{H}{def\tilde{\kappa}}$$

$$(12) \quad \tilde{\tau}_2 = \frac{3}{2} \cdot \frac{Hb}{f\Delta_{ebd}}$$

Observe that (11) and (12) both use the minimal number of points (i.e. four) for approximating τ . One might argue that this induces an asymmetry in the numerical results. However, this is something easy to fix in a practical situation (for example by averaging the results obtained in the two directions).

We justify (11) as follows: It is known that $\tau = \frac{-\alpha_s \times \alpha_{ss} \cdot \alpha_{sss}}{\kappa^2}$, [2]. In other words, $|\tau|$ is the component of α_{sss} which is perpendicular to the osculating plane. We start by writing down a finite difference scheme for α_{sss} (for justifications, see [5] §5.4).

$$(13) \quad \alpha_{sss}(0) \approx 6 \cdot \frac{\frac{\alpha(\epsilon_2) - \alpha(\epsilon_1)}{\epsilon_2 - \epsilon_1} - \frac{\alpha(\epsilon_1) - \alpha(0)}{\epsilon_1} - \frac{\alpha(\epsilon_1) - \alpha(0)}{\epsilon_1} - \frac{\alpha(0) - \alpha(-\delta)}{\delta}}{\delta + \epsilon_2}.$$

When ϵ_1 and δ approach zero, the plane defined by P_{i-1}, P_i and P_{i+1} approaches the osculating plane. More precisely, if δ and ϵ_1 are small enough, we find (from the local canonical form of the curve) that the equation of the plane passing through $\alpha(-\delta), \alpha(0)$ and $\alpha(\epsilon)$ can be written as

$$Ax + By + z = 0$$

with $A = \frac{\kappa\tau}{6}\epsilon_1\delta + O(3)$ and $B = -\frac{\tau}{3}(\epsilon_1 - \delta) + O(2)$. Here, x, y and z are the Frenet frame coordinates. The normal vector to this plane is $\frac{(A, B, 1)}{\sqrt{A^2 + B^2 + 1}}$ which is a first order approximation of $(0, 0, 1)$, the normal vector to the osculating plane. We have

$$\begin{aligned} \kappa\tau &= \alpha_{sss} \cdot (0, 0, 1) \\ &= \alpha_{sss} \cdot \frac{(A, B, 1)}{\sqrt{A^2 + B^2 + 1}} + O(1). \end{aligned}$$

But

$$\alpha(-\delta) \cdot (A, B, 1) = \alpha(0) \cdot (A, B, 1) = \alpha(\epsilon_1) \cdot (A, B, 1) = 0$$

and

$$\alpha(\epsilon_2) \cdot \frac{(A, B, 1)}{\sqrt{A^2 + B^2 + 1}} = \pm H.$$

Therefore

$$\begin{aligned} \kappa\tau &\approx \frac{\pm 6H}{(\epsilon_2 - \epsilon_1)\epsilon_2(\delta + \epsilon_2)} \\ &\approx \pm \frac{6H}{def}. \end{aligned}$$

From the previous section we know that $\kappa \approx \pm \tilde{\kappa} = \frac{4\Delta_{abc}}{abc}$ and therefore $\kappa \approx \pm \frac{6H}{def\tilde{\kappa}}$.

There is an interesting similarity between (11) and $\tilde{\kappa}$. This is easily seen if we write $\tilde{\kappa}$ in a slightly different way. By basic geometry argument we can say that $\tilde{\kappa} = \pm \frac{2h}{ab}$ where h is the height of the triangle with sides a , b and c . Thus $\tilde{\kappa}\tau$ can be looked at as a three dimensional version of $\tilde{\kappa}$. This is one reason why we like (11). Another reason is that it provides us with an easy visual understanding of the torsion similar to the understanding we already have of the curvature. It is indeed easier to evaluate the height and sides of a tetrahedron together with the radius of the osculating circle than a derivative. It makes the evaluation of the torsion from a picture of the curve very intuitive.

For finding (12), we started from the following definition: if $P = \alpha(t)$ and $Q = \alpha(0)$, then $\tau = \lim_{P \rightarrow Q} \frac{\theta}{t}$, where θ is the angle between the osculating planes at P and Q respectively [4]. We approximated θ by $\sin \tilde{\theta}$ where $\tilde{\theta}$ denotes the angle between the plane $P_i P_{i+1} P_{i+2}$ and the plane $P_{i-1} P_i P_{i+1}$. If we assign to the plane $P_{i-1} P_i P_{i+1}$ the equation $z = 0$ and if we let the segment $P_i P_{i+1}$ represent the x-axis, then the plane passing through BCD has equation $Hy - D_y z = 0$ where D_y is the projection of the point D on the y-axis.

Given the equations of two planes $A_1x + B_1y + C_1z + D_1 = 0$ and $A_2x + B_2y + C_2z + D_2 = 0$, the angle between the two planes is given by

$$\cos \theta = \frac{A_1A_2 + B_1B_2 + C_1C_2}{\sqrt{A_1^2 + B_1^2 + C_1^2} \sqrt{A_2^2 + B_2^2 + C_2^2}}.$$

A straightforward computation gives us

$$\begin{aligned} \sin \tilde{\theta} &= \pm \frac{H}{\sqrt{H^2 + D_y^2}} \\ &= \pm \frac{H}{\sqrt{e^2 - D_x^2}} \\ &= \pm \frac{H}{\sqrt{e^2 - (\frac{e^2 + b^2 - d^2}{2b})^2}} \\ &= \pm \frac{2Hb}{\sqrt{4e^2b^2 - (e^2 + b^2 - d^2)^2}} \\ &= \pm \frac{2Hb}{4\Delta_{ebd}} \\ &= \pm \frac{Hb}{2\Delta_{ebd}}. \end{aligned} \tag{14}$$

On the other hand, we can use the local canonical form to approximate the equations of the plane $P_{i-1}P_iP_{i+1}$ and $P_iP_{i+1}P_{i+2}$ in the Frenet frame coordinates. Respectively, we have approximately

$$\left(\frac{\kappa\tau\epsilon_1\delta}{6}\right)x + \left(\frac{\tau}{3}\right)(\epsilon_1 - \delta)y + z = 0$$

and

$$\left(\frac{-\kappa\tau\epsilon_1\epsilon_2}{6}\right)x + \left(\frac{\tau}{3}\right)(\epsilon_1 + \epsilon_2)y + z = 0.$$

So

$$\cos^2 \tilde{\theta} \approx \frac{(1 + (\frac{\tau}{3})^2(\epsilon_1 - \delta)(\epsilon_1 + \epsilon_2))^2}{1 + (\frac{\tau}{3}(\epsilon_1 - \delta))^2 + (\frac{\tau}{3}(\epsilon_1 + \epsilon_2))^2}$$

and therefore

$$\sin^2 \tilde{\theta} \approx \frac{\tau^2}{9}(\epsilon_2 + \delta)^2 \approx \frac{\tau^2}{9}f^2$$

i.e. $\frac{3\sin\tilde{\theta}}{f} \approx \pm\tau$. Combining this result with (14), we get (12).

With the help of the symbolic computation software Mathematica and using the local canonical form of a curve, we computed the Taylor series expansion for both expressions. For (11), we obtained

$$(15) \quad \tilde{\tau}_1 = \tau + \frac{\tau\kappa_s}{6\kappa}(a - b + 3e) + \frac{\tau_s}{4}(b - a + e) + O(2).$$

For (12) we obtained

$$(16) \quad \tilde{\tau}_2 = \tau + \frac{\tau\kappa_s}{6\kappa}(a + b + e) + \frac{\tau_s}{4}(b - a + e) + O(2).$$

These expansions prove that (11) and (12) are not equivalent.

There are many ways to approximate τ_s . For the reasons mentioned above, we chose to use (11). In a similar way as for κ_s , we used (15) to obtain the following numerically invariant expression for τ_s . For symmetry reasons and in view of the results obtained for κ_s , we decided on using a centered formula. The result is the following five point approximation.

$$(17) \quad \tilde{\tau}_s(P_i) = 4 \cdot \frac{\tilde{\tau}^1(P_{i+1}) - \tilde{\tau}^1(P_{i-1}) + (2a + 2b - 2d - 3h + g) \frac{\tilde{\tau}^1(P_i)\tilde{\kappa}_s(P_i)}{6\tilde{\kappa}(P_i)}}{2a + 2b + 2d + h + g}$$

Where $g = d(P_{i-2}, P_{i-1})$ as illustrated on figure 1 and $\tilde{\kappa}_s$ can be taken to be either (4), (5) or (6).

3.3. Numerical test. We considered the curve $\alpha(t) = (\cos t, \sin t, \sqrt{t})$ whose graph is given in figure 12. Note that α is not parametrized by arc length. The signature

of this curve is given by the four following quantities.

$$\begin{aligned}\kappa &= |\alpha_{ss}| = \frac{|\alpha_t \times \alpha_{tt}|}{|\alpha_t|^3} = \frac{2\sqrt{16t^3 + 4t^2 + 1}}{(1 + 4t)^{3/2}}, \\ \kappa_s &= \frac{d\kappa}{dt} \frac{dt}{ds} = \frac{\frac{d\kappa}{dt}}{\frac{ds}{dt}} = \frac{8(8t^2 + 2t - 3)\sqrt{t}}{\sqrt{16t^3 + 4t^2 + 1}(1 + 4t)^3}, \\ \tau &= -\frac{\alpha_t \times \alpha_{tt} \cdot \alpha_{ttt}}{|\alpha_t \times \alpha_{tt}|^2} = \frac{-2\sqrt{t}(3 + 4t^2)}{16t^3 + 4t^2 + 1}, \\ \tau_s &= \frac{d\tau}{dt} \frac{dt}{ds} = \frac{\frac{d\tau}{dt}}{\frac{ds}{dt}} = \frac{2(64t^5 - 16t^4 + 240t^3 + 16t^2 - 3)}{\sqrt{4t + 1}(16t^3 + 4t^2 + 1)^2}.\end{aligned}$$

We tested our numerical expressions on the portion of the curve given by $\frac{\pi}{2} \leq t \leq \frac{3\pi}{2}$ for different partition size. We used $\tilde{\kappa}_s$ given by equation (6) for evaluating $\tilde{\tau}_s$. Care was taken to partition the curve very irregularly. In the example presented on figures 13-22, we computed the signature for a partition $\{t_i\}_{i=0}^N$ built the following way

$$\begin{aligned}t_0 &= o \\ t_1 &= t_0 + \Delta t \\ t_2 &= t_1 + \frac{\Delta t}{2} \\ t_3 &= t_2 + \frac{\Delta t}{3} \\ t_4 &= t_3 + \Delta t \\ t_5 &= t_4 + \frac{\Delta t}{2} \\ t_6 &= t_5 + \frac{\Delta t}{3} \text{ etc.}\end{aligned}$$

The resulting graphs are presented for $\Delta t = 0.1, 0.05, 0.025$ and 0.0125 . For simplicity, we graphed two projections of the signature curve. Although a lot of the information is lost this way, we believe that it gives a good measure of the effectivity of the method. We can see that although the partition is not regular, the graphs seem to converge to the exact graphs when Δt becomes small. Therefore we believe that our formulas work properly for generic small partitions.

4. CONCLUSION

We now have in our hands good approximations for the Euclidean and affine signature of a planar curve. The formulas we obtained are invariant under the action of the Euclidean and affine group respectively. Another important property is that they are valid for any fine partition of the given curve. In a near future, we hope to use them for the recognition of planar curves and perhaps also for some other applications of the signature curve among the numerous possible ones. In particular, we wish to improve the numerical results obtained with noisy images and prove that the signature can be a practical tool of object recognition.

We also have good approximations for the four differential invariants which parametrize the Euclidean signature of a space curve, namely the curvature, the torsion and their derivatives with respect to arc length. The formulas we obtained are Euclidean invariants and valid for any fine partition of the given curve, although

better results are obtained with equidistant partitions. In particular, we showed in section 3 that $\tilde{\tau} = \pm 6 \cdot \frac{H}{def\kappa}$ is a good approximation for the torsion. We believe that this expression has a value on its own as it gives an easy visual understanding of the concept of torsion. Space curve recognition (for example blood vessels and trajectory of particles) is an interesting computer vision problem where our formulas could find applications.

FIGURE 2. Exact Data Used for Testing the Planar Euclidean Case.

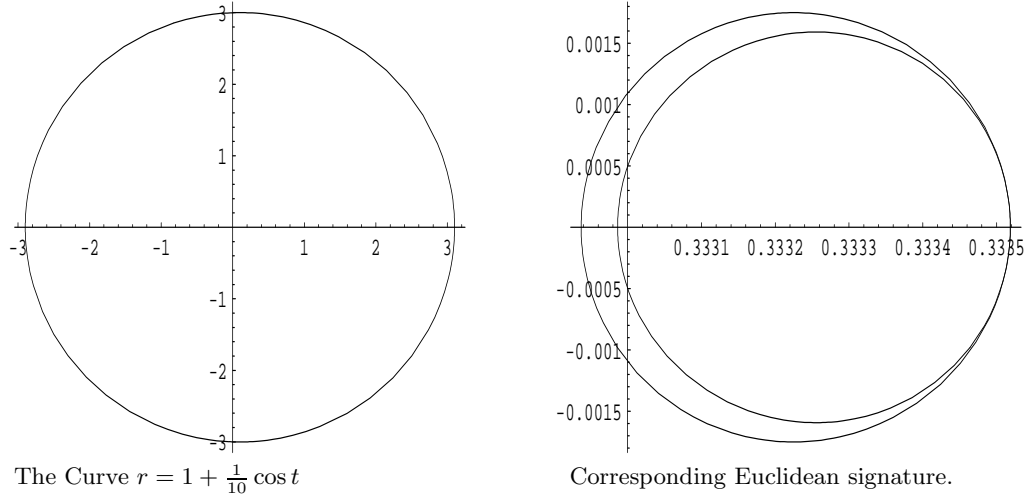


FIGURE 3. Three Different Partitions of the Initial Curve Used in the Planar Euclidean Case.

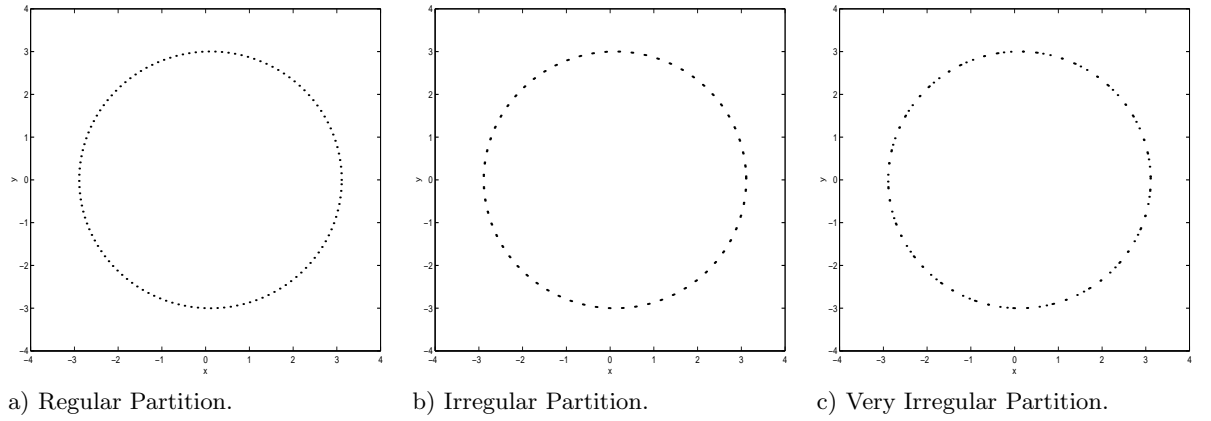
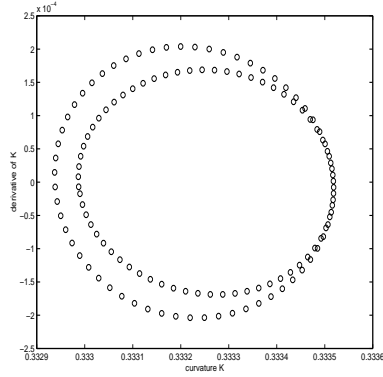
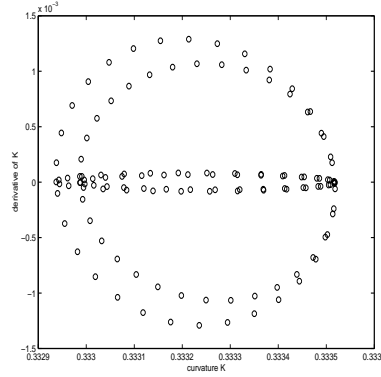


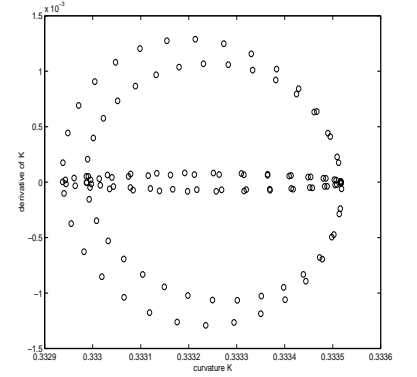
FIGURE 4. Approximations of the Euclidean Signature Curve Obtained with (2).



With Partition 3-a.

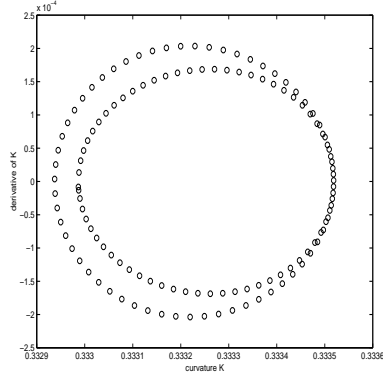


With Partition 3-b.

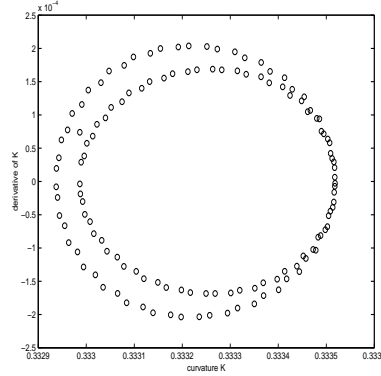


With Partition 3-c.

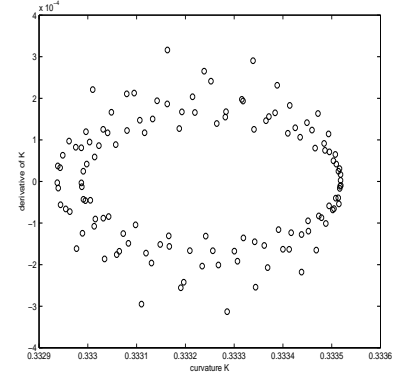
FIGURE 5. Approximations of the Euclidean Signature Curve Obtained with (3).



With Partition 3-a.



With Partition 3-b.



With Partition 3-c.

FIGURE 6. Approximations of the Euclidean Signature Curve Obtained with (4).

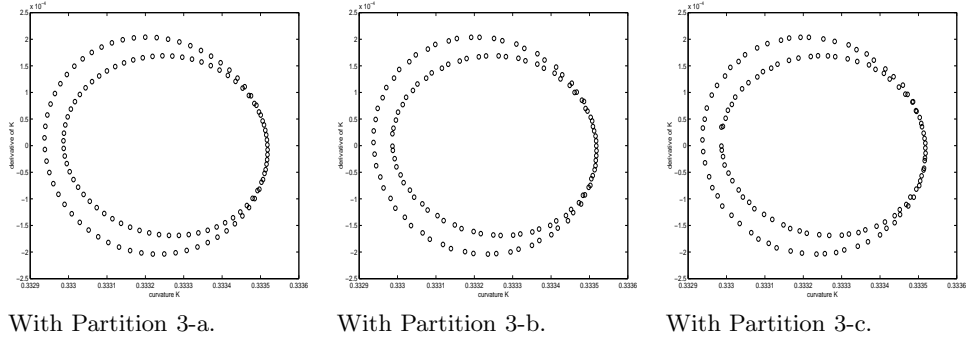


FIGURE 7. Approximations of the Euclidean Signature Curve Obtained with (5).

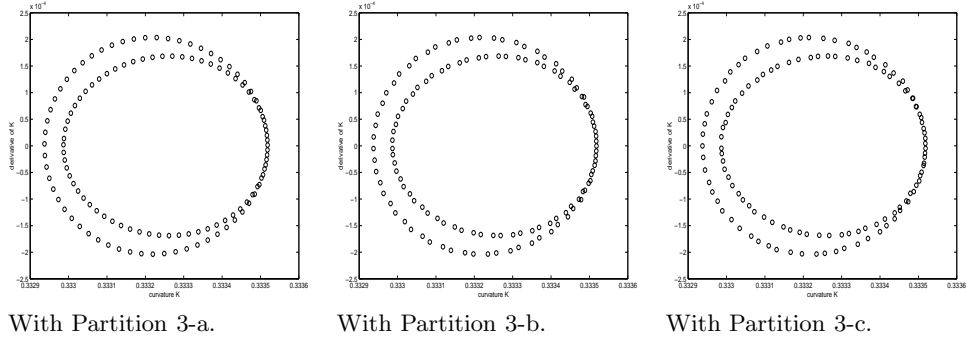


FIGURE 8. Approximation of the Euclidean Signature Curve Obtained with (6).

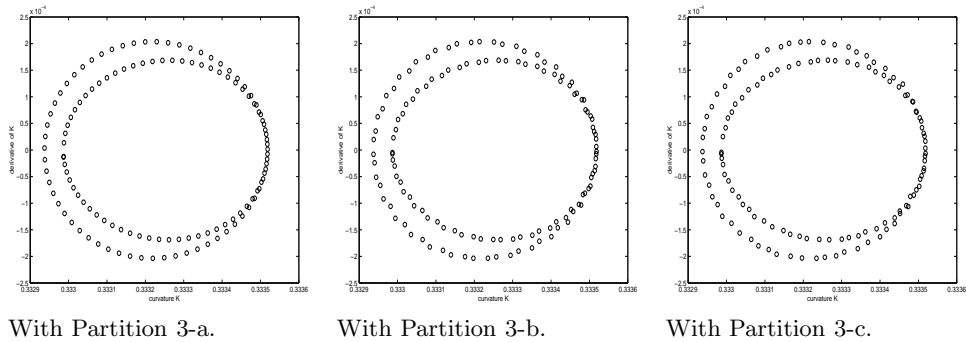


FIGURE 9. Exact Data Used for Testing the Planar Affine Case.

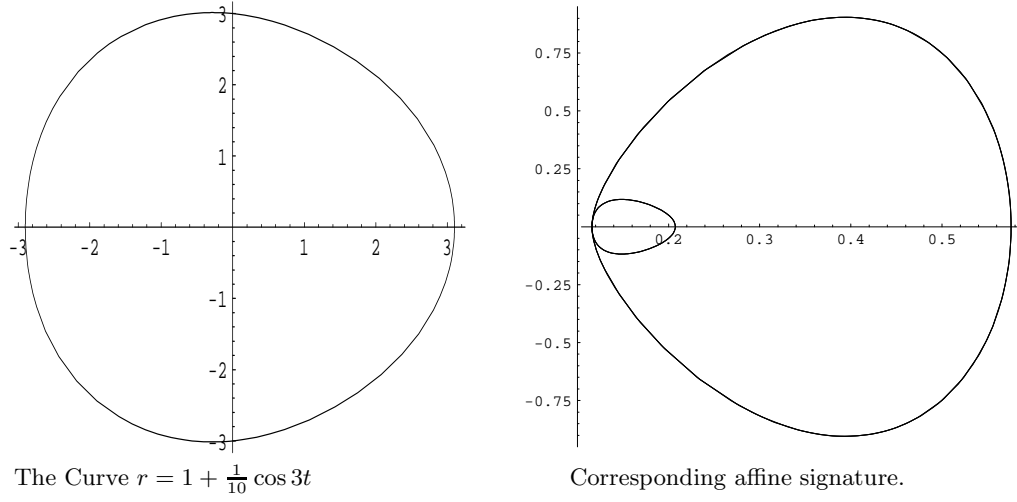


FIGURE 10. Three Different Partitions of the Initial Curve Used in the Planar Affine Case.

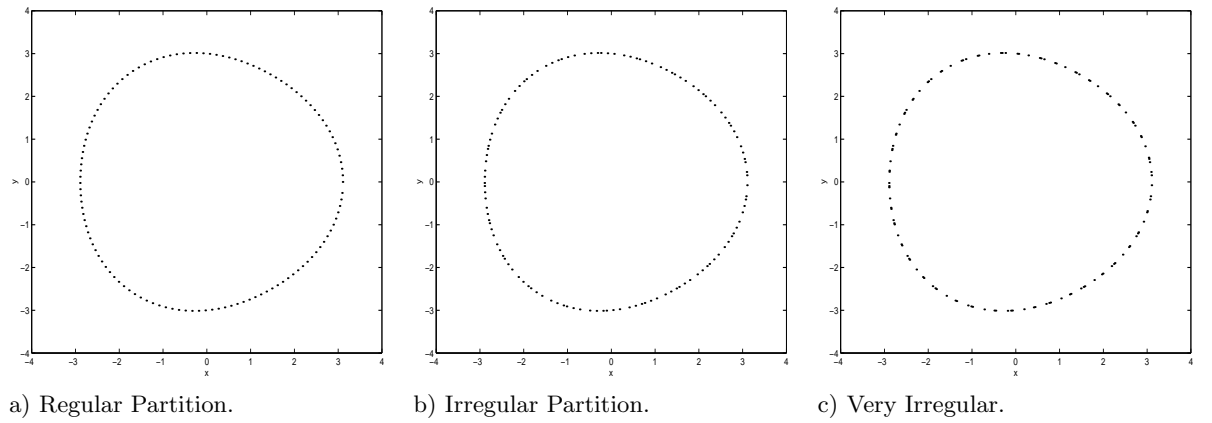
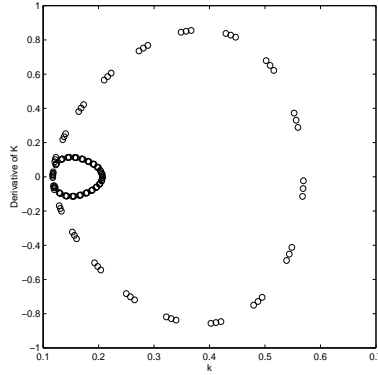
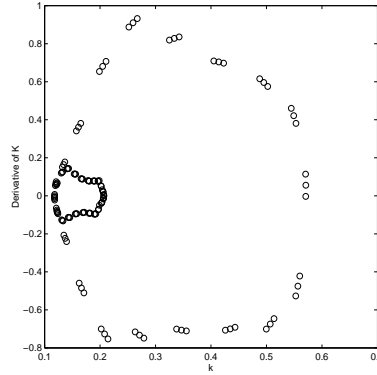


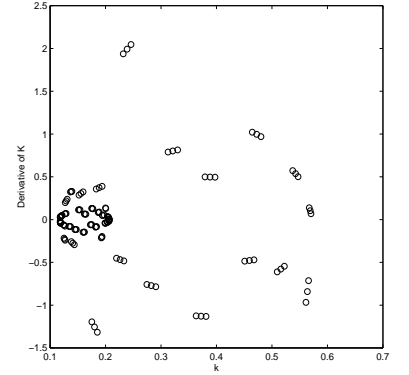
FIGURE 11. Approximations of the Affine Signature Curve Obtained with (9).



With partition 10-a.

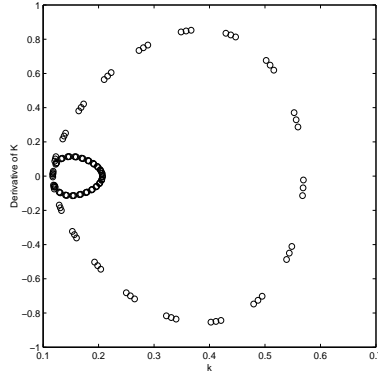


With partition 10-b.

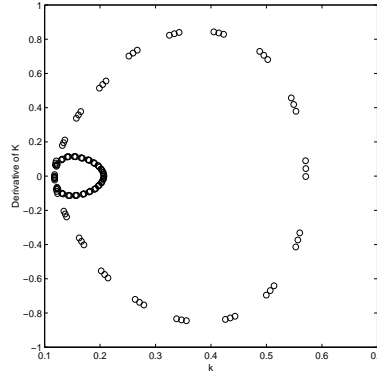


With partition 10-c.

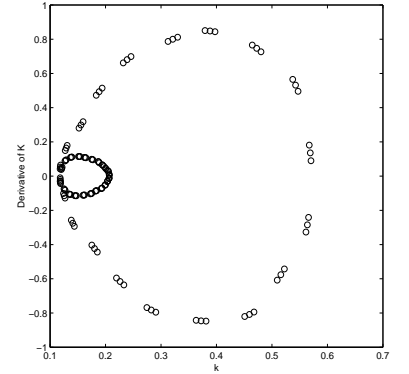
FIGURE 12. Approximations of the Affine Signature Curve Obtained with (10).



With Partition 10-a.

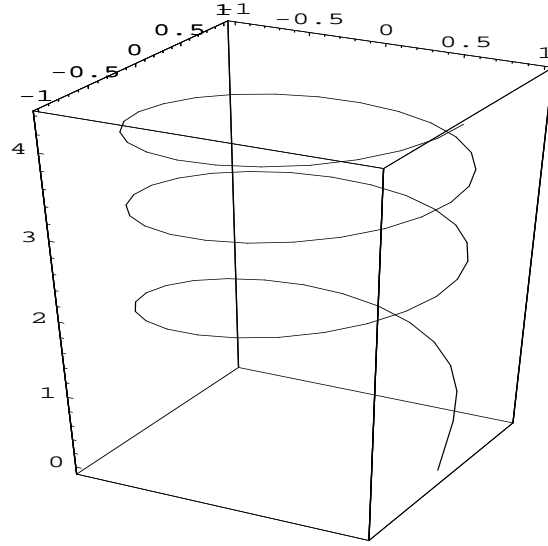


With Partition 10-b.



With Partition 10-c.

FIGURE 13. Curve Used to Test the Formulas Proposed for the Spatial Signature



The Curve $\alpha(t) = (\cos t, \sin t, \sqrt{t})$, for $0 \leq t \leq 6\pi$

FIGURE 14. Projections of the Exact Signature for the Curve Above

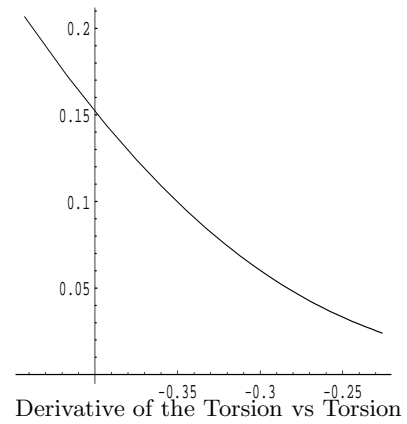
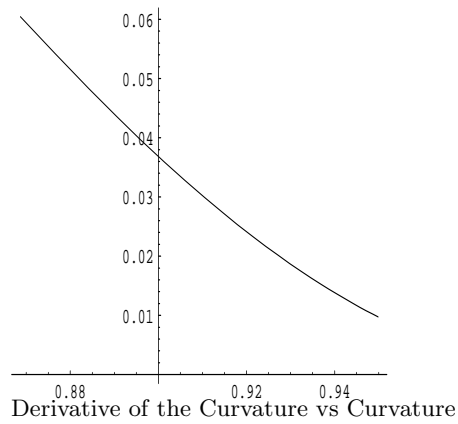


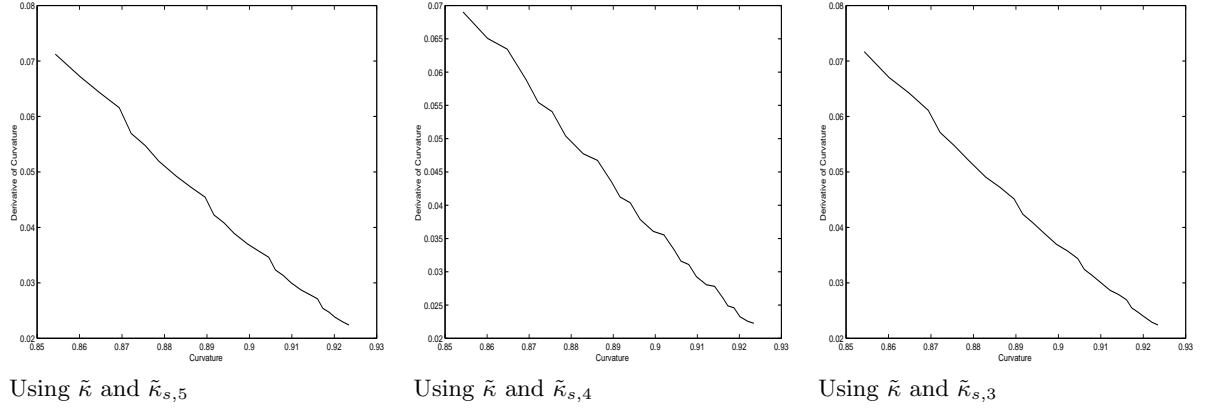
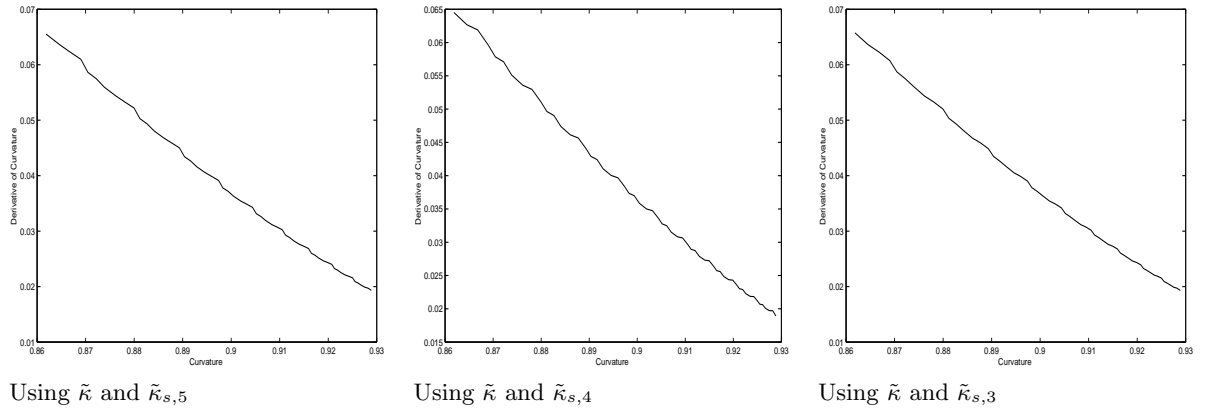
FIGURE 15. Approximations Obtained with $\Delta t = 0.1$

 FIGURE 16. Approximations Obtained with $\Delta t = 0.05$


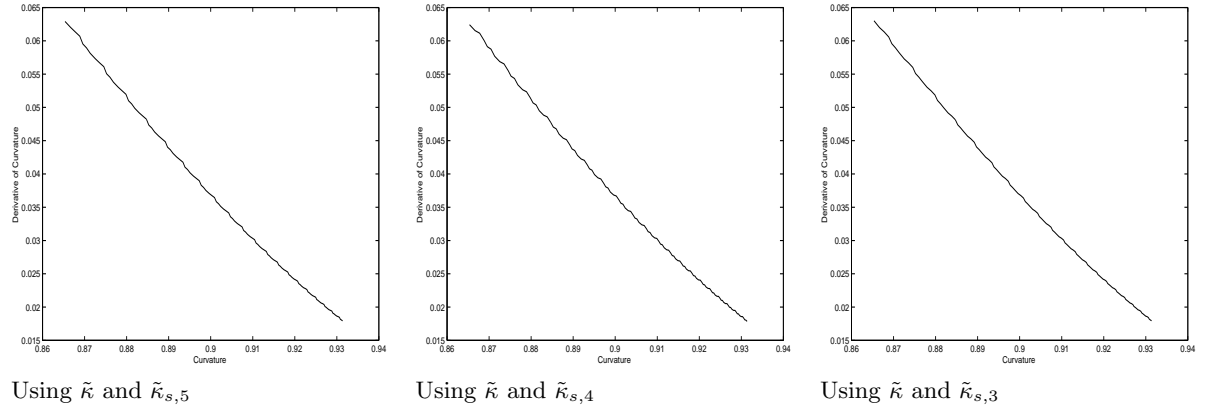
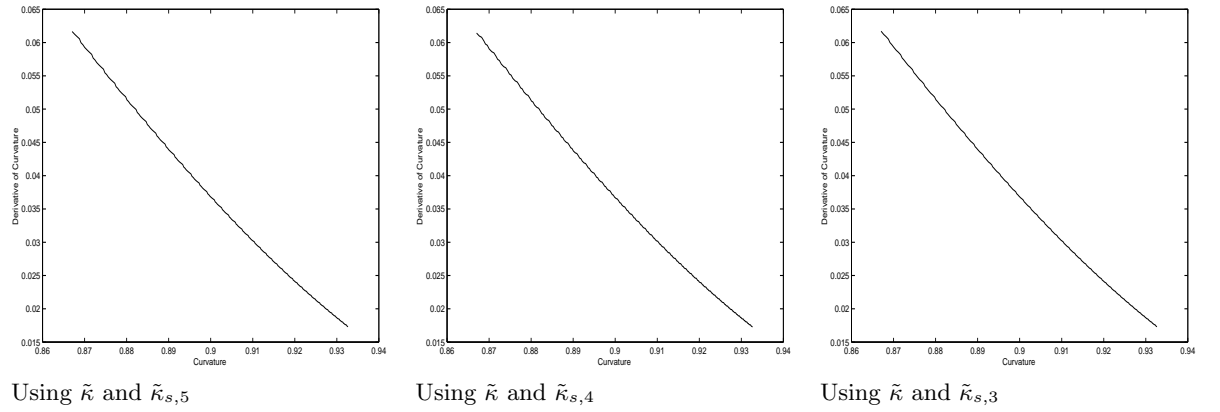
FIGURE 17. Approximations Obtained with $\Delta t = 0.025$ FIGURE 18. Approximations Obtained with $\Delta t = 0.0125$ 

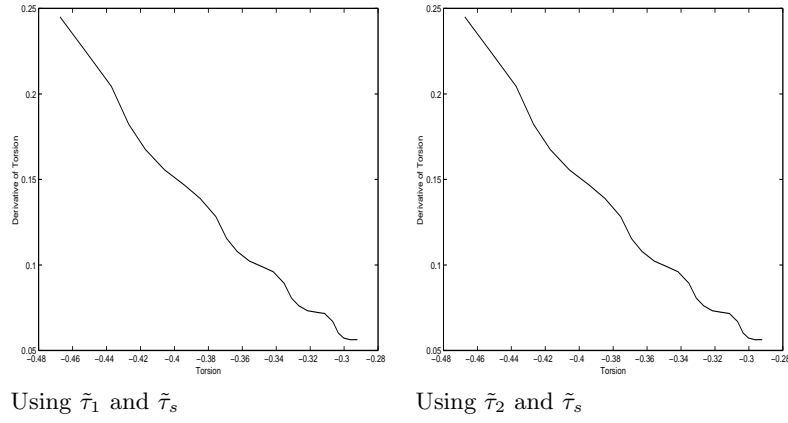
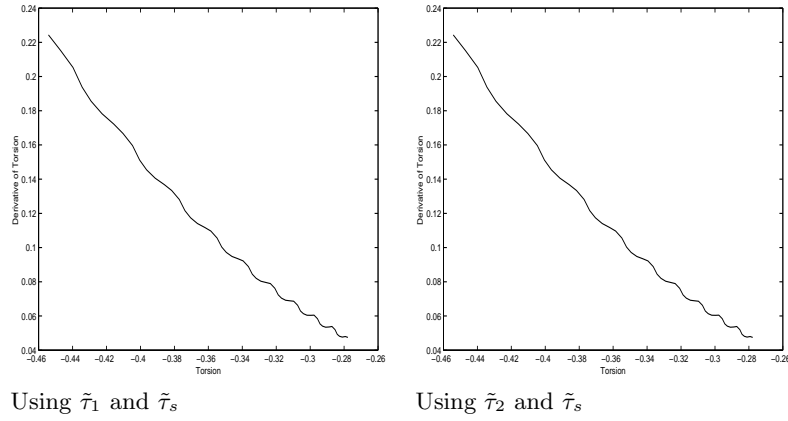
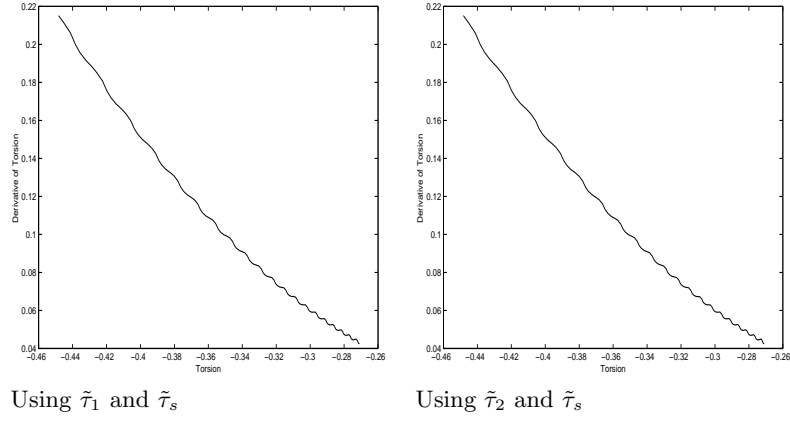
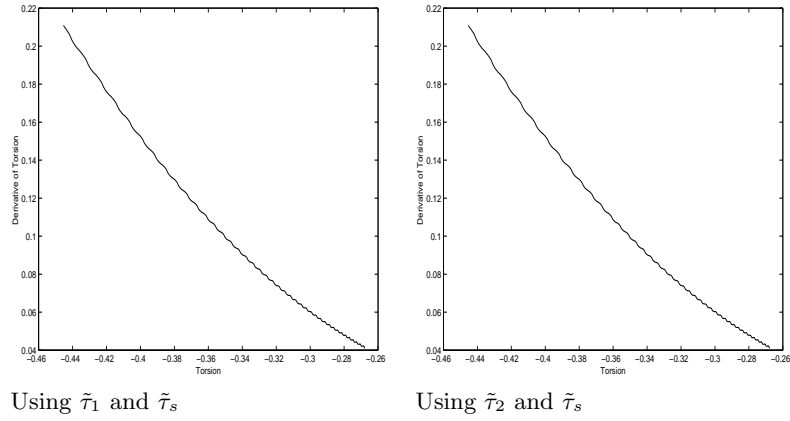
FIGURE 19. Approximations Obtained with $\Delta t = 0.1$

 FIGURE 20. Approximations Obtained with $\Delta t = 0.05$


FIGURE 21. Approximations Obtained with $\Delta t = 0.025$ FIGURE 22. Approximations Obtained with $\Delta t = 0.0125$ 

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